

# Relation between Distributional and Leray-Hopf Solutions to the Navier-Stokes Equations

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## Preliminary Considerations

Cauchy problem for the Navier-Stokes equations:

$$\left. \begin{aligned} \partial_t v + v \cdot \nabla v &= \Delta v - \nabla p \\ \operatorname{div} v &= 0 \end{aligned} \right\} \text{ in } \mathbb{R}^3 \times (0, \infty)$$
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$$v(x, 0) = v_0(x) \quad x \in \mathbb{R}^3.$$

**Leray-Hopf solutions:** For any  $v_0 \in L^2_\sigma(\mathbb{R}^3)$  there is

$$v \in L^\infty(0, T; L^2_\sigma) \cap L^2(0, T; H^1), \quad \text{all } T > 0$$

$$\lim_{t \rightarrow 0^+} \|v(t) - v_0\|_2 = 0$$

$$\int_{\mathbb{R}^3} |v(x, t)|^2 + 2 \int_0^t \int_{\mathbb{R}^3} |\nabla v(x, \tau)|^2 \leq \int_{\mathbb{R}^3} |v_0(x)|^2,$$

all  $t \in [0, T]$ .

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**Energy Equality** (Prodi-Lions)

$$v \in L^{\frac{2s}{s-3}}(0, T; L^s), \quad s > 3, \quad (\text{B}),$$

or

$$v \in C([0, T]; L^3), \quad (\text{C})$$

**Uniqueness/Regularity** (Prodi-Serrin-Ladyzhenskaya)

# Main Results

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Then  $v$  is necessarily a Leray-Hopf solution.

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In fact, if **merely**  $v \in \{(A), (B), (C)\}$ , both total kinetic energy and dissipation

$$\int_{\mathbb{R}^3} |v(x, t)|^2, \quad \int_0^t \int_{\mathbb{R}^3} |\nabla v(x, \tau)|^2,$$

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are **infinite at all times**.

Our result implies, in particular, that this cannot happen if only the kinetic energy of the initial datum is finite.

# Main Results: Precise Statement

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**Definition.** Let  $v_0 \in L^2_\sigma$ . Then  $v \in L^2_{\text{loc}}(\mathbb{R}^3 \times [0, T))$  is a *distributional solution* (corresponding to  $v_0$ ) if

$$\int_0^T \int_{\mathbb{R}^3} v \cdot (\partial_t \varphi + \Delta \varphi + v \cdot \nabla \varphi) = - \int_{\mathbb{R}^3} v_0 \cdot \varphi(0)$$

$$\int_0^T \int_{\mathbb{R}^3} v \cdot \nabla \phi = 0,$$

for all

$$\varphi \in \mathcal{D}_T := \{\varphi \in C_0^\infty(\mathbb{R}^3 \times [0, T)) : \operatorname{div} \varphi = 0\},$$

and  $\phi \in C_0^\infty(\mathbb{R}^3 \times (0, T))$ .

**Theorem 1.** *Let  $v$  be a distributional solution satisfying one of the following conditions:*

$$v \in L^4(0, T; L^4);$$

$$v \in L^{\frac{2s}{s-3}}(0, T; L^s), s > 3;$$

$$v \in C([0, T]; L^3).$$

*Then  $v$  is a Leray-Hopf solution.*

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**Theorem 2.** *Let  $v$  be a distributional solution satisfying one of the following conditions:*

$$v \in L^{\frac{2s}{s-3}}(\delta, T; L^s), s > 3;$$

$$v \in C([\delta, T]; L^3),$$

*all “small”  $\delta > 0$ , and also*

$$\lim_{t \rightarrow 0^+} \int_{\mathbb{R}^3} (v(x, t) - v_0(x)) \cdot \Phi(x) dx = 0, \quad \text{all } \Phi \in L^2_\sigma.$$



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**Corollary 2.** *If  $v$  is a distributional solution satisfying one of the conditions*

$$v \in L^{\frac{2s}{s-3}}(0, T; L^s), s > 3; \quad v \in C([0, T]; L^3),$$

*then  $v$  is the unique Leray-Hopf solution corresponding to  $v_0$ , and also  $v \in C^\infty((0, T] \times \mathbb{R}^3)$ .*

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**Corollary 3.** *If  $v$  is a distributional solution such that*

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**Corollary 4.** *If  $v$  is a distributional solution satisfying one of the conditions for all "small"  $\delta > 0$*

$$v \in L^{\frac{2s}{s-3}}(\delta, T; L^s), \quad s > 3; \quad v \in C([0, \delta, T]; L^3),$$

*and  $v(t) \rightarrow v_0$  weakly in  $L^2_\sigma$  as  $t \rightarrow 0^+$ , then  $v$  is a Leray-Hopf solution, so that  $v \in C^\infty((0, T] \times \mathbb{R}^3)$ .*

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### 2. Investigation in either class

$$v \in L^{\frac{2s}{s-3}}(0, T; L^s), \quad s > 3; \quad v \in C([0, T]; L^3),$$

falls, instead, in the framework of “very weak” (or “mild”) solutions, and has been performed extensively.

## Some Remarks

FOIAS (1961):

$$v \in L^2(0, T; H^1) \cap L^r(0, T; L^s), s > 3, r > \frac{2s}{s-3}; v_0 \in L^2_\sigma.$$

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KATO (1984), GIGA (1986):

$$v \in C([0, T]; L^3) \cap L^{\frac{2s}{s-3}}(0, T; L^s), s > 3, v_0 \in L^2_\sigma \cap L^3.$$

$$\lim_{t \rightarrow 0} \left( t^{\frac{s-3}{2s}} \|v(t)\|_s \right) = 0.$$

**3.** BUCKMASTER & VICOL (2017/2018), LUO (2018) show non-uniqueness in the class of distributional solutions defined by

$$v \in C([0, T]; H^\beta), \beta > 0 \text{ "small"}, \text{ and } v_0 \in L^2_\sigma.$$

having, therefore, infinite dissipation.



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Our results imply, in particular, that any possible non-uniqueness result in that class could only hold with  $\beta < 1/2$ .

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$$v \in L^\infty(0, T; L^3)$$

continues to ensure that the distributional solution  $v$  is in the Leray-Hopf class.

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**5.** Our results continue to hold in (1) bounded or (2) exterior domain (no-slip conditions), and (3) in a torus (periodic boundary conditions). The proof in cases (1), (2) becomes much more technical.

# Basic Ideas

**Step 1.** For the **given** distributional solution  $v$ , consider the **linear** Cauchy problem

$$\left. \begin{aligned} \partial_t u + v \cdot \nabla u &= \Delta u - \nabla p \\ \operatorname{div} u &= 0 \end{aligned} \right\} \text{ in } \mathbb{R}^3 \times (0, T)$$
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$$u(x, 0) = v_0(x) \quad x \in \mathbb{R}^3.$$

By standard Galerkin method one shows the existence of a solution

$$u \in L^\infty(0, T; L^2_\sigma) \cap L^2(0, T; H^1),$$
$$u(t) \rightarrow v_0 \quad \text{in } L^2.$$

## Basic Ideas

Set

$$w := v - u$$

then

$$\int_0^T \int_{\mathbb{R}^3} w \cdot (\partial_t \varphi + \Delta \varphi + v \cdot \nabla \varphi) = 0$$
$$\int_0^T \int_{\mathbb{R}^3} w \cdot \nabla \phi = 0,$$

for all

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In which case, we get

$$\int_0^T \int_{\mathbb{R}^3} w \cdot f = 0, \quad \text{for all } f \in C_0^\infty(\mathbb{R}^3 \times (0, T)),$$

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implying that  $v \equiv u$  is a Leray-Hopf solution.

This step can be achieved if  $v$  is in one of the classes (A), (B) or (C).

Implementation of Step 2, under the assumption

$$v \in L^r(0, T; L^s), \quad r := \frac{2s}{s-3}, \quad s > 3.$$

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“Maximal regularity” class ( $p, q \in (1, \infty)$ ):

$$\mathcal{W}^{p,q} := \{ \psi \in W^{1,p}(0, T; L^q_\sigma) \cap L^p(0, T; W^{2,q}) \}.$$

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**Lemma 1.** (Technical) *Test space  $\mathcal{D}_T$  is dense in*

$$\left\{ \psi \in \mathcal{W}^{p_1, q_1} \cap \mathcal{W}^{p_2, q_2} : \psi(T) = 0 \right\},$$

*for any  $p_i, q_i \in (1, \infty)$ ,  $i = 1, 2$ .*

Recall that  $(r := 2s/(s - 3))$

$$w := v - u \in L^r(0, T; L^s) + L^\infty(0, T; L^2_\sigma).$$

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**Lemma 2.** (Embeddings + Hölder) *The linear form*

$$\psi \in \mathcal{W}^{1,2} \cap \mathcal{W}^{r',s'} \rightarrow \int_0^T \int_{\mathbb{R}^3} v \cdot \nabla \psi \cdot w \in \mathbb{R}$$

*is continuous* ( $r' := r/(r - 1)$ ,  $s' := s/(s - 1)$ .)

## Basic Ideas

As a consequence,

$$\int_0^T \int_{\mathbb{R}^3} w \cdot (\partial_t \varphi + \Delta \varphi + v \cdot \nabla \varphi) = 0$$

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implies

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implies

$$\int_0^T \int_{\mathbb{R}^3} w \cdot \underbrace{(\partial_t \psi + \Delta \psi + v \cdot \nabla \psi)}_{= f \in C_0^\infty(\mathbb{R}^3 \times (0, T))} = 0$$

for all

$$\psi \in W^{1,2} \cap \mathcal{W}^{r',s'}, \text{ with } \psi(T) = 0.$$

**Lemma 3.** *Let*

$$\alpha \in L^r(0, T; L^s_\sigma) \equiv L^{r,s}, \quad r := \frac{2s}{s-3}.$$

*For any  $F \in C_0^\infty(\mathbb{R}^3 \times (0, T))$ , the problem*

$$\left. \begin{aligned} \partial_t \Psi - \alpha \cdot \nabla \Psi &= \Delta \Psi - \nabla p + F \\ \operatorname{div} \Psi &= 0 \end{aligned} \right\} \text{ in } \mathbb{R}^3 \times (0, T),$$
$$\Psi(x, 0) = 0 \quad x \in \mathbb{R}^3,$$

*has at least one solution  $\Psi \in \mathcal{W}^{1,2} \cap \mathcal{W}^{r',s'}$ .*

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Estimate in  $\mathcal{W}^{r',s'}$ . For  $\lambda > 0$ , set

$$\zeta = e^{-\lambda t} \Psi, \quad \rho = e^{-\lambda t} p,$$

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$$G = e^{-\lambda t} F + \alpha \cdot \nabla \zeta.$$

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We get the Stokes problem

$$\left. \begin{aligned} \partial_t \zeta + \lambda \zeta &= \Delta \zeta - \nabla \rho + G \\ \operatorname{div} \zeta &= 0 \end{aligned} \right\} \text{in } \mathbb{R}^3 \times (0, T),$$
$$\zeta(x, 0) = 0 \quad x \in \mathbb{R}^3,$$

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Maximal regularity ( $\lambda$  “large”):

$$\begin{aligned} \|\zeta\|_{\mathcal{W}^{r',s'}} + \lambda \|\zeta\|_{L^{r',s'}} &\leq c \|G\|_{L^{r',s'}} \\ &\leq c (\|\alpha \cdot \nabla \zeta\|_{L^{r',s'}} + \|F\|_{L^{r',s'}}) . \end{aligned}$$

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Split  $\alpha = \alpha_1 + \alpha_2$  with

$$\|\alpha_1\|_{L^{r,s}} < \varepsilon, \quad \alpha_2 \in L^{\infty,\infty} .$$

We get (Hölder + embeddings + Ehrling)

$$\begin{aligned}\|\alpha \cdot \nabla \zeta\|_{L^{r',s'}} &\leq \varepsilon \|\nabla \zeta\|_{L^{\frac{r}{r-2}, \frac{s}{s-2}}} + \|\alpha_2\|_{L^{\infty,\infty}} \|\nabla \zeta\|_{L^{r',s'}} \\ &\leq c\varepsilon \|\zeta\|_{\mathcal{W}^{r',s'}} + c_\varepsilon \|\alpha_2\|_{L^{\infty,\infty}} \|\zeta\|_{L^{r',s'}} .\end{aligned}$$

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Therefore:

$$\|\zeta\|_{\mathcal{W}^{r',s'}} \leq c \|F\|_{L^{r',s'}} .$$



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Test the original equation with  $\Psi$ ,  $\Delta\Psi$ , and  $\partial_t\Psi$ :

$$\frac{1}{2} \frac{d}{dt} \|\Psi\|_2^2 + \|\nabla\Psi\|_2^2 = \int_{\mathbb{R}^3} F \cdot \Psi$$

$$\frac{1}{2} \frac{d}{dt} \|\nabla\Psi\|_2^2 + \|\Delta\Psi\|_2^2 = \int_{\mathbb{R}^3} (\alpha \cdot \nabla\Psi - F) \cdot \Delta\Psi$$

$$\frac{1}{2} \frac{d}{dt} \|\nabla\Psi\|_2^2 + \|\partial_t\Psi\|_2^2 = - \int_{\mathbb{R}^3} (\alpha \cdot \nabla\Psi - F) \cdot \partial_t\Psi$$

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Assumption on  $\alpha$  + Sobolev lead to

$$\|\Psi\|_{\mathcal{W}^{1,2}} \leq c \|F\|_{L^{2,2}} .$$

**QED**

Proof of Leray-Hopf property.

## Proof of Leray-Hopf property.

For a given  $f \in C_0^\infty(\mathbb{R}^3 \times (0, T))$  we use the previous lemma with

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Thus,  $\Psi(T - t) \equiv \psi(t) \in \mathcal{W}^{1,2} \cap \mathcal{W}^{r',s'}$ , and solves the final-value problem

$$\left. \begin{aligned} \partial_t \psi + v \cdot \nabla \psi + \Delta \psi &= \nabla p + f \\ \operatorname{div} \psi &= 0 \end{aligned} \right\} \text{ in } \mathbb{R}^3 \times (0, T),$$
$$\psi(x, T) = 0 \quad x \in \mathbb{R}^3,$$

## Basic Ideas

Recalling that  $w := v - u$  satisfies

$$\int_0^T \int_{\mathbb{R}^3} w \cdot (\partial_t \psi + \Delta \psi + v \cdot \nabla \psi) = 0$$

for all

$$\psi \in W^{1,2} \cap \mathcal{W}^{r',s'}, \text{ with } \psi(T) = 0,$$

we get

$$\int_0^T \int_{\mathbb{R}^3} (v - u) \cdot f = 0, \text{ for all } f \in C_0^\infty(\mathbb{R}^3 \times (0, T))$$

that is  $v \equiv u$ .

**QED**

## Basic Ideas

Sketch of the proof under the alternate assumption

$$v \in L^{\frac{2s}{s-3}}(\delta, T; L^s), \quad s > 3; \quad v \rightarrow v_0 \text{ weakly in } L^2.$$



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$$v \in L^{\frac{2s}{s-3}}(\delta, T; L^s), \quad s > 3; \quad v \rightarrow v_0 \text{ weakly in } L^2.$$

Let

$$\chi_\delta(t) = \begin{cases} 1 & \text{if } t \geq 2\delta \\ 0 & \text{if } t \leq \delta \end{cases}, \quad |\chi'_\delta| \leq \frac{c}{\delta}.$$

## Basic Ideas

Sketch of the proof under the alternate assumption

$$v \in L^{\frac{2s}{s-3}}(\delta, T; L^s), \quad s > 3; \quad v \rightarrow v_0 \text{ weakly in } L^2.$$

Let

$$\chi_\delta(t) = \begin{cases} 1 & \text{if } t \geq 2\delta \\ 0 & \text{if } t \leq \delta \end{cases}, \quad |\chi'_\delta| \leq \frac{c}{\delta}.$$

One can show ( $w := v - u$ )

$$\int_\delta^T \chi_\delta \int_{\mathbb{R}^3} w \cdot (\partial_t \psi + \Delta \psi + v \cdot \nabla \psi) = - \int_\delta^{2\delta} \chi'_\delta \int_{\mathbb{R}^3} w \cdot \psi,$$

for all  $\psi \in W^{1,2} \cap \mathcal{W}^{r',s'}$ ,  $\psi(T) = 0$ .

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$$\int_{\delta}^T \chi_{\delta} \int_{\mathbb{R}^3} w \cdot f = - \int_{\delta}^{2\delta} \chi'_{\delta} \int_{\mathbb{R}^3} w \cdot \psi, \quad f \in C_0^{\infty}(\mathbb{R}^3 \times (0, T)).$$

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Letting  $\delta \rightarrow 0$ ,

$$\begin{aligned} \int_0^T \int_{\mathbb{R}^3} w \cdot f &= - \lim_{\delta \rightarrow 0} \int_{\delta}^{2\delta} \chi'_{\delta} \int_{\mathbb{R}^3} w \cdot \psi \\ &= 0 \end{aligned}$$

because

$$\chi'_{\delta} \sim \delta^{-1}, \quad w \rightarrow 0 \text{ weakly in } L^2, \quad \psi \in C([0, T]; L^2).$$

**QED**

Entirely similar arguments can be employed to draw the same conclusions if either

$$v \in L^4(0, T; L^4) \quad [\text{resp. } v \in L^4(\delta, T; L^4)],$$

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with the mentioned conditions on the attainability of the initial datum  $v_0$ .

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Thank you